

能带理论(下)

(I) 晶格电子的平均速度

① 布洛赫波 ψ_{nk} , 计算平均速度.

$$\left\{ \begin{array}{l} \vec{V}_n(k) = \frac{1}{m} \int \psi_{nk}^* \hat{\vec{p}} \psi_{nk} d\vec{r} \\ \hat{\vec{p}} = -i\hbar \vec{\nabla} \end{array} \right.$$

代入布洛赫波形式 $\psi_{nk} = e^{i\vec{k}\cdot\vec{r}} \underline{\underline{u}_{nk}(\vec{r})}$

$$\begin{aligned} \Rightarrow \vec{V}_n(k) &= \frac{1}{m} \int e^{-i\vec{k}\cdot\vec{r}} \underline{\underline{u}_{nk}^*(\vec{r})} (-i\hbar \vec{\nabla}) [e^{i\vec{k}\cdot\vec{r}} \underline{\underline{u}_{nk}(\vec{r})}] \cdot d\vec{r} \\ &= \frac{1}{m} \int e^{-i\vec{k}\cdot\vec{r}} \underline{\underline{u}_{nk}^*(\vec{r})} [\hbar \vec{k} e^{i\vec{k}\cdot\vec{r}} \underline{\underline{u}_{nk}} + e^{i\vec{k}\cdot\vec{r}} \cdot (-i\hbar \vec{\nabla} \underline{\underline{u}_{nk}})] \\ &= \frac{1}{m} \int \underline{\underline{u}_{nk}^*(\vec{r})} (\hbar \vec{k} - i\hbar \vec{\nabla}) \underline{\underline{u}_{nk}(\vec{r})} d\vec{r} \\ &= \frac{1}{m} \int \underline{\underline{u}_{nk}^*(\vec{r})} (\hbar \vec{k} + \hat{\vec{p}}) \underline{\underline{u}_{nk}(\vec{r})} d\vec{r} \\ \vec{V}_n(k) &= \frac{1}{m} \langle \underline{\underline{u}_{nk}} | \underline{\underline{\hat{p}}} | \underline{\underline{u}_{nk}} \rangle \quad \text{周期解} \end{aligned}$$

② $u_{nk}(r)$ to Schrödinger Equation

$$\left[\frac{(\vec{p} + \hbar \vec{k})^2}{2m} + V(r) \right] u_{nk}(\vec{r}) = E_n(k) u_{nk}(\vec{r})$$

$$H_K = \frac{(\vec{p} + \hbar \vec{k})^2}{2m} + V(r),$$

$$H_K \cdot u_{nk}(\vec{r}) = E_n(k) u_{nk}(\vec{r})$$

$$\text{左边 } D_K: \quad \text{右边} = (\vec{p}_K H_K) u_{nk}(r) + t_K (\vec{p}_K u_{nk}(r))$$

$$\hbar \cdot \frac{(\vec{t} \vec{k} + \vec{p})}{m} U_{nk}(\vec{r}) + H_k \vec{D}_k U_{nk}(\vec{r})$$

$$\text{左边} = \vec{D}_k [\Sigma_n(k) U_{nk}(r)]$$

$$= \underbrace{\vec{D}_k \Sigma_n(k) U_{nk}(r)} + \Sigma_n(k) \vec{D}_k U_{nk}(r)$$

$$\frac{\hbar}{m} (\vec{t} \vec{k} + \vec{p}) U_{nk}(\vec{r}) + H_k \vec{D}_k U_{nk}(\vec{r}) = \underline{\vec{D}_k \Sigma_n(k) U_{nk}} + \underline{\Sigma_n(k) \vec{D}_k U_{nk}}$$

左侧 $U_{nk}(r)$, 并不约分 (Dirac notation)

$$\begin{aligned} & \frac{\hbar}{m} \langle U_{nk} | \vec{t} \vec{k} + \vec{p} | U_{nk} \rangle + \langle U_{nk} | H_k | \vec{D}_k U_{nk} \rangle \\ &= \cancel{\langle U_{nk} | U_{nk} \rangle} \cdot \vec{D}_k \Sigma_n(k) + \Sigma_n(k) \cancel{\langle U_{nk} | \vec{D}_k U_{nk} \rangle} \end{aligned}$$

$$\Rightarrow \frac{1}{m} \langle \vec{t} \vec{k} + \vec{p} \rangle = \frac{1}{\hbar} \vec{D}_k \Sigma_n(k)$$

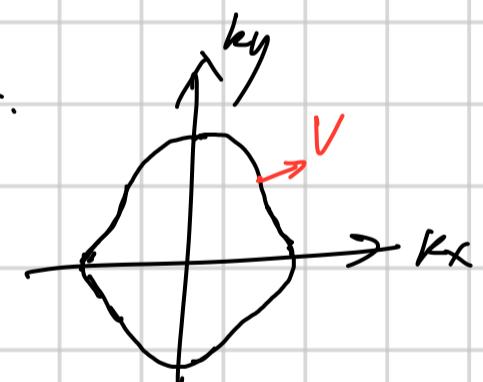
得到 $\vec{V}_n(k) = \frac{1}{\hbar} \vec{D}_k \Sigma_n(k)$ 晶格电子平均速度.

③ 讨论 $\Sigma \sim \hbar \omega$, $V_n \sim D_k \omega$ 波包速度.

- 沿导角面梯度方向

- 自由电子气 $\Sigma(k) = \frac{t^2 k^2}{2m}$

$$V(k) = \frac{1}{\hbar} \nabla_k \left(\frac{t^2 k^2}{2m} \right) = \vec{t} \vec{k} / m$$



II. 电子运动的半经典模型 / 运动方程.

- 晶格电子对外场的响应 (\vec{E}, \vec{B}) \Rightarrow 经典处理
- 晶格电子的运动状态 \Rightarrow 是 力 (能带理论)

(1) \vec{F} 外力 (下), 布洛赫电子运动方程.

电子能量的改变 = 外力做功

$$\Delta \Sigma(k) = \vec{f} \cdot \Delta \vec{s} \quad \frac{d\Sigma(k)}{dt} = \vec{f} \cdot \vec{v}$$

$$\text{左边: } \frac{d\Sigma(k)}{dt} = \vec{\nabla}_k \Sigma(k) \cdot \frac{dk}{dt} = \hbar \vec{v} \cdot \frac{dk}{dt} = \frac{d(\hbar k)}{dt} \cdot \vec{v}$$

$$\text{右边: } \vec{f} \cdot \vec{v}$$

$$\Rightarrow \frac{d(\hbar k)}{dt} = \vec{f} \quad \text{运动方程.}$$

讨论: ① \vec{f} 仅包含外力 (不包含晶格势能)

电磁场中: $\frac{d(\hbar k)}{dt} = -e \left[\underbrace{\vec{E}(\vec{r}, t)}_{\parallel f} + \underbrace{\vec{V}_n(\vec{k}) \times \vec{B}(\vec{r}, t)}_{\perp f} \right]$

② 牛顿运动方程 (经典)
 $\hbar k$ 描述电子状态 (量子)

(2) 加速度与有效质量

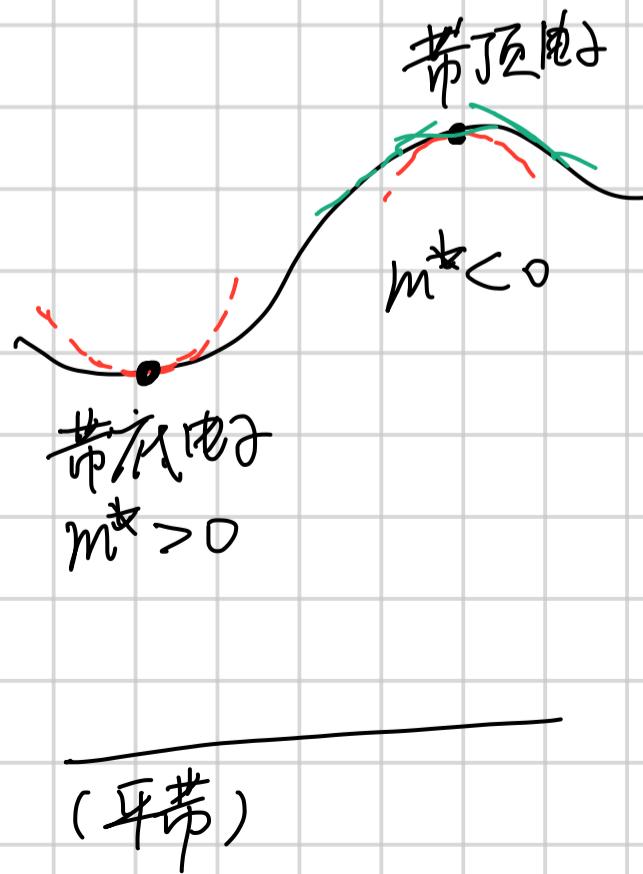
$$\begin{aligned} \vec{a} &= \frac{d\vec{v}}{dt} = \frac{d}{dt} \left[\frac{1}{\hbar} \vec{\nabla}_k \Sigma_n(k) \right] \\ &= \frac{1}{\hbar} \vec{\nabla}_k \left[\vec{\nabla}_k \Sigma_n(k) \right] \cdot \frac{dk}{dt} \\ &\quad \boxed{C \quad f = m^* a} \end{aligned}$$

$$\Rightarrow \text{有效质量 } m^* = \frac{t^2}{\nabla_k^2 \Sigma_n(k)}$$

$$\text{讨论: ① 1D, } m^* = t^2 \cdot \frac{1}{d^2 \Sigma_n(k)/dk^2}$$

$$\text{平坦能带 } m^* = t^2 \cdot \frac{1}{d^2/ds^2 \rightarrow 0} \rightarrow \infty$$

\Rightarrow 局域 (局域化)



② 3D 情况

$$\vec{a} = \frac{1}{\hbar^2} \vec{\nabla}_k [\vec{P}_k \ln(k)] \cdot \vec{f}$$

$$\begin{pmatrix} a_x \\ a_y \\ a_z \end{pmatrix} = \frac{1}{m^*} \begin{pmatrix} f_x \\ f_y \\ f_z \end{pmatrix}, \quad \frac{1}{m^*} = \frac{1}{\hbar^2} \frac{\partial^2 \Sigma_h(h)}{\partial k_x \partial k_p}, \quad (\alpha, \beta = x, y, z)$$

矩阵形式.

$$\frac{1}{m^*} = \frac{1}{\hbar^2} \begin{pmatrix} \frac{\partial^2 \Sigma}{\partial k_x^2} & \frac{\partial^2 \Sigma}{\partial k_x \partial k_y} & \frac{\partial^2 \Sigma}{\partial k_x \partial k_z} \\ \frac{\partial^2 \Sigma}{\partial k_y \partial k_x} & \frac{\partial^2 \Sigma}{\partial k_y^2} & \frac{\partial^2 \Sigma}{\partial k_y \partial k_z} \\ \frac{\partial^2 \Sigma}{\partial k_z \partial k_x} & \frac{\partial^2 \Sigma}{\partial k_z \partial k_y} & \frac{\partial^2 \Sigma}{\partial k_z^2} \end{pmatrix}$$

✓ (对角化 $\frac{1}{m^*}$ 矩阵)
找主轴方向

✓ 对于不平行主轴的外力, 产生不平行于主轴的加速度。
(滑板电子不仅反向受外力, 还有滑板内势能)

III. 稳恒电场中的电子运动

① 运动方程 $t \frac{d\vec{k}}{dt} = -e \vec{E}$ ← 不会时

$$t_k(t) - t_k(0) = -e \vec{E} \cdot t$$

$$\rightarrow \vec{k}(t) = \vec{k}(0) - e \vec{E} \cdot t / t_0 \quad \checkmark$$

稳恒电场下, 电子沿矢量移动。

② 运动速度分析 (半正, 半反)

$$\Sigma(\vec{k}) = \Sigma(-\vec{k}) \Rightarrow v(\vec{k}) = -v(-\vec{k})$$

(\vec{k} 和 $-\vec{k}$ 处运动速度相反)

✓ 无矢量, 对称单元 \Rightarrow 无电流
(满或不满)

